

Constraints on Warm Dark Matter models from high-redshift long gamma-ray bursts

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ABSTRACT

Structures in Warm Dark Matter (WDM) models are exponentially suppressed below a certain scale, characterized by the dark matter particle mass, m_x . Since structures form hierarchically, the presence of collapsed objects at high-redshifts can set strong lower limits on m_x . We place robust constraints on m_x using recent results from the *Swift* database of high-redshift gamma-ray bursts (GRBs). We parameterize the redshift evolution of the ratio between the cosmic GRB rate and star formation rate (SFR) as $\propto (1+z)^\alpha$, thereby allowing astrophysical uncertainties to partially mimic the cosmological suppression of structures in WDM models. Using a maximum likelihood estimator on two different $z > 4$ GRB subsamples (including two bursts at $z > 8$), we constrain $m_x \gtrsim 1.6\text{--}1.8$ keV at 95% CL, when marginalized over a flat prior in α . We further estimate that 5 years of a SVOM-like mission would tighten these constraints to $m_x \gtrsim 2.3$ keV. Our results show that GRBs are a powerful probe of high-redshift structures, providing robust and competitive constraints on m_x .

Key words: methods: statistical–gamma-ray burst: general–cosmology: dark matter

1 INTRODUCTION

The current concordance cosmology, in which structure formation proceeds in a hierarchal manner driven by pressureless cold dark matter (CDM), has been remarkably successful in explaining the observed properties of large-scale structures in the Universe (e.g., Tegmark et al. 2006; Benson 2010) and the cosmic microwave background (CMB) (e.g., Komatsu et al. 2011). Such observables probe scales in the range ~ 1 Gpc down to ~ 10 Mpc. On smaller scales, $\lesssim 1$ Mpc, there are still some discrepancies between standard Λ CDM and observations (e.g., Menci et al. 2012). For instance, N-body simulations predict more satellite galaxies than are observed both around our galaxy (the so-called “missing satellite problem”; Moore et al. e.g., 1999; Klypin et al. e.g., 1999), and in the field as recently noted by the ALFALFA survey (e.g., Papastergis et al. 2011; Ferrero et al. 2012). Furthermore, simulations of the most massive Galactic CDM subhaloes are too centrally condensed to be consistent with the kinematic data of the bright Milky Way satellites (e.g., Boylan-Kolchin et al. 2011).

Moreover, observations of small galaxies show that their central density profile is shallower than predicted by CDM N-body simulations (e.g., Moore 1994; de Blok et al. 2001; Donato et al. 2009; Macciò et al. 2012; Governato et al. 2012).

Baryonic feedback is a popular prescription for resolving such discrepancies. Feedback caused by supernovae (SNe) explosions and heating due to the UV background may suppress the baryonic content of low-mass haloes (e.g., Governato et al. 2007; Mashchenko et al. 2008; Busha et al. 2010; Sobacchi & Mesinger 2013b), and make their inner density profile shallower (e.g., de Souza & Ishida 2010; de Souza et al. 2011). However, accurately matching observations is still difficult even when tuning feedback recipes (e.g. Boylan-Kolchin et al. 2012).

An alternative explanation might be found if dark matter (DM) consisted of lower mass (\sim keV) particles and thus was “warm” (WDM; e.g., Bode et al. 2001; Khlopov & Kouvaris 2008; de Vega & Sanchez 2012; de Vega et al. 2012; Kang et al. 2013; Destri et al. 2013; Kamada et al. 2013). The resulting effective pressure and free-streaming would decrease structure on small-scales, though again fine-tuning might be required to fully

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match all the observations (e.g. Boylan-Kolchin et al. 2011; Macciò et al. 2012; Borriello et al. 2012).

The most powerful testbed for these scenarios is the high-redshift Universe. Structure formation in WDM models (or in any cosmological model with an equivalent power-spectrum cut-off) is exponentially suppressed on small-scales (e.g., Schneider et al. 2012, 2013). Since structures form hierarchically, these small halos are expected to host the first galaxies. If indeed dark matter were sufficiently “warm”, the high-redshift Universe would be empty. Therefore, the mere presence of a galaxy at high-redshift can set strong lower limits on the WDM particle mass.

Due to their high luminosity, gamma-ray bursts (GRBs) constitute a remarkable tool to probe the high- z Universe and small-scale structures. They provide a glimpse of the first generations of stars (e.g., de Souza et al. 2011, 2012), as well as provide constraints on primordial non-Gaussianity (Maio et al. 2012). As pointed out by Mesinger et al. (2005), the detection of a single GRB at $z > 10$ would provide very strong constraints on WDM models.

Here we extend the work of Mesinger et al. (2005) by *presenting robust lower limits on WDM particle masses, using the latest Swift GRB data*. The current data, including many redshift measurements, allows us to perform an improved statistical analysis by directly comparing the distribution of bursts in various models as a function of redshift. Furthermore, we make more conservative¹ assumptions throughout the analysis, such as normalizing the SFR-GRB ratio at high redshifts (thereby using a shorter, more accurate lever arm which minimizes modeling uncertainty), using an unbiased luminosity function and allowing the SFR-GRB ratio to evolve with redshift. Finally, we study the effectiveness of future observations in improving the current constraints.

Current limits on dark matter masses, m_x , are motivated by several observations. The Lyman- α forest implies $m_x \gtrsim 1$ keV (e.g., Viel et al. 2008) and $m_{\nu_s} > 8$ keV for sterile neutrinos (Seljak et al. 2006; Boyarsky et al. 2009). Likewise, WDM models with a too warm candidate ($m_x < 0.75$ keV) cannot simultaneously reproduce the stellar mass function and the Tully-Fisher relation (Kang et al. 2013). Also, the fact that reionization occurred at $z \gtrsim 6$ implies $m_x \gtrsim 0.75$ keV (Barkana et al. 2001). However, all of these limits are strongly affected by a degeneracy between astrophysical (i.e. baryonic) processes and the dark matter mass. Our approach in this work is more robust, driven only by the shape of the redshift evolution of the $z > 4$ SFR. Furthermore, it is important to note that the SFR is exponentially attenuated at high-redshifts in WDM models. Since astrophysical uncertainties are unable to mimic such a rapid suppression, probes at high-redshifts (such as GRBs and reionization) are powerful in constraining WDM cosmologies.

The outline of this paper is as follows. In §2 we discuss how we derive the dark matter halo mass function and SFR in WDM and CDM models. In §3 we derive the correspond-

ing GRB redshift distribution. In §4 we discuss the adopted observed GRB sample. In §5 we present our analysis and main results. In §6, we discuss possible future constraints using a theoretical mock sample. Finally, in §7, we present our conclusions.²

2 STRUCTURE FORMATION IN A WDM DOMINATED UNIVERSE

Massive neutrinos from the standard model (SM) of particle physics were one of the first dark matter candidates. However, structures formed in this paradigm are incompatible with observations. Other alternative dark matter candidates usually imply an extension of the SM. The DM particle candidates span several order of magnitude in mass (Boyarsky et al. 2009): axions with a mass of $\sim 10^{-6}$ eV, first introduced to solve the problem of CP violation in particle physics, supersymmetric (SUSY) particles (gravitinos, neutralinos, axinos) with mass in the range \sim eV-GeV, superheavy dark matter, also called *Wimpzillas*, (also considered as a possible solution to the problem of cosmic rays observed above the GZK cutoff), Q-balls, and sterile neutrinos with mass \sim keV range, just to cite a few. For a review about dark matter candidates see Bertone et al. (2005). Two promising candidates for warm dark matter are the sterile neutrino (Dodelson & Widrow 1994; Shaposhnikov & Tkachev 2006) and gravitino (Ellis et al. 1984; Moroi et al. 1993; Kawasaki et al. 1997; Primack 2003; Gorbunov et al. 2008).

In WDM models the growth of density perturbations is suppressed on scales smaller than the free streaming length. The lighter the WDM particle, the larger the scale below which the power spectrum is suppressed. In addition to this power-spectrum cutoff, one must also consider the residual particle velocities. As described in Barkana et al. (2001), these act as an effective pressure, slowing the early growth of perturbations. Below we describe how we include these two effects in our analysis (for more information, please see Barkana et al. 2001; Mesinger et al. 2005).

2.1 Power spectrum cutoff

The free-streaming scale, below which the linear perturbation amplitude is suppressed (e.g., Colombi et al. 1996; Bode et al. 2001; Viel et al. 2005), is given by the comoving scale

$$R_{fs} \approx 0.11 \left(\frac{\Omega_x h^2}{0.15} \right)^{1/3} \left(\frac{m_x}{\text{keV}} \right)^{-4/3} \text{Mpc}, \quad (1)$$

where Ω_x is the total energy density contained in WDM particles relative to the critical density, h is the Hubble constant in units of $100 \text{ km s}^{-1} \text{ Mpc}^{-1}$, and m_x is the WDM particle mass.

The resulting modification of the matter power spectrum can be computed by multiplying the CDM power

¹ Throughout the text, we use “conservative” to imply biasing the GRB distribution towards lower redshifts. This is conservative since it mimics the effects of WDM, thereby resulting in weaker constraints on the particle mass.

² Throughout the paper, we adopt the cosmological parameters from WMAP9 (Hinshaw et al. 2012), $\Omega_m = 0.264$, $\Omega_\Lambda = 0.736$, $n_s = 0.97$, $\sigma_8 = 0.8$ and $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

spectrum $P_{\text{CDM}}(k)$ by an additional transfer function (Bode et al. 2001):

$$P_{\text{WDM}}(k) = P_{\text{CDM}}(k) [1 + (\epsilon k)^{2\mu}]^{-5\mu}, \quad (2)$$

where $\mu = 1.12$ and

$$\epsilon = 0.049 \left(\frac{\Omega_x}{0.25}\right)^{0.11} \left(\frac{m_x}{\text{keV}}\right)^{-1.11} \left(\frac{h}{0.7}\right)^{1.22} h^{-1} \text{Mpc}. \quad (3)$$

2.2 Effective pressure

Structure formation in WDM models will be further suppressed by the residual velocity dispersion of the WDM particles, which delays the growth of perturbations. As shown in Barkana et al. (2001) and Mesinger et al. (2005), this “effective pressure” is of comparable importance to the power spectrum cutoff in determining the WDM mass functions. The pressure can be incorporated in the halo mass function by raising the critical linear extrapolated overdensity threshold at collapse, $\delta_c(M, z)$. Using spherically symmetric hydrodynamics simulations, and exploiting the analogy between the WDM effective pressure and gas pressure, Barkana et al. (2001) computed the modified $\delta_c(M, z)$. They showed that one can define an effective WDM Jeans mass, i.e., the mass scale below which collapse is significantly delayed by the pressure. We will denote this scale as:

$$M_{\text{WDM}} \approx 1.8 \times 10^{10} \left(\frac{\Omega_x h^2}{0.15}\right)^{1/2} \left(\frac{m_x}{\text{keV}}\right)^{-4}, \quad (4)$$

where we make the standard assumption of a fermionic spin 1/2 particle. Note that equation (4) differs from the original one proposed by Barkana et al. (2001) by a factor of 60. With this adjustment, we find that the collapsed fractions, $F_{\text{coll}}(z)$, simulated assuming a sharp mass cutoff at M_{WDM} are in good agreement with the full random-walk procedure (assuming a more gradual rise in $\delta_c(M)$) of Barkana et al. (2001) (see Fig. 2). The former approach has the advantage of being considerably faster and simpler.

2.3 Collapsed fraction of haloes and cosmic star formation

We use the Sheth-Tormen mass function, f_{ST} , (Sheth & Tormen 1999) to estimate the number density of dark matter haloes, $n_{\text{ST}}(M, z)$, with mass greater than M :

$$f_{\text{ST}} = A \sqrt{\frac{2a_1}{\pi}} \left[1 + \left(\frac{\sigma^2}{a_1 \delta_c^2}\right)^p\right] \frac{\delta_c}{\sigma} \exp\left[-\frac{a_1 \delta_c^2}{2\sigma^2}\right], \quad (5)$$

where $A = 0.3222$, $a_1 = 0.707$, $p = 0.3$ and $\delta_c = 1.686$. The mass function f_{ST} can be related to $n_{\text{ST}}(M, z)$ by

$$f_{\text{ST}} = \frac{M}{\rho_m} \frac{dn_{\text{ST}}(M, z)}{d \ln \sigma^{-1}}, \quad (6)$$

where ρ_m is the total mass density of the background Universe. The variance of the linearly extrapolated density field $\sigma(M, z)$ is given by

$$\sigma^2(M, z) = \frac{b^2(z)}{2\pi^2} \int_0^\infty k^2 P(k) W^2(k, M) dk, \quad (7)$$

where $b(z)$ is the growth factor of linear perturbations normalized to $b = 1$ at the present epoch, and $W(k, M)$ is a real space top-hat filter. In order to calculate the CDM power spectrum $P_{\text{CDM}}(k)$, we use the CAMB code³. For WDM models, we compute the power spectrum, $P_{\text{WDM}}(k)$, with equation (2).

The fraction of mass inside collapsed haloes, $F_{\text{coll}}(> M_{\text{min}}, z)$, is then given by:

$$F_{\text{coll}}(z) = \frac{1}{\rho_m} \int_{M_{\text{min}}}^\infty dM M n_{\text{ST}}(M, z), \quad (8)$$

and the minimum mass is estimated as

$$M_{\text{min}} = \text{Max}[M_{\text{gal}}(z), M_{\text{WDM}}(m_x)], \quad (9)$$

where M_{gal} corresponds to the minimum halo mass capable of hosting star forming galaxies. We use $M_{\text{gal}} \equiv \bar{M}_{\text{min}}$ from equation (11) in Sobacchi & Mesinger (2013a), who present a physically-motivated expression for the evolution of M_{gal} which takes into account a gas cooling criterion as well as radiative feedback from an ionizing UV background (UVB) during inhomogeneous reionization⁴. In other words, M_{gal} is set by astrophysics, whereas $M_{\text{WDM}}(m_x)$ is set by the particle properties of dark matter. In Fig. 1, we show these two limits as a function of redshift. As we shall see, our conclusions are derived from the regime where $M_{\text{WDM}} > M_{\text{gal}}$, and hence they are not sensitive to the exact value of M_{gal} . Another point worth highlighting is that models with $m_x \gtrsim 3.5$ keV have $M_{\text{WDM}}(m_x) \geq M_{\text{gal}}$ up to $z \approx 11$. Thus, any constraint beyond ~ 3.5 keV will no longer be robust, since galaxy formation can be suppressed even in CDM, below halo masses comparable to M_{WDM} . Therefore, observations at much higher redshifts are required to obtain tighter constraints.

In figure 2, we plot the fraction of the total mass collapsed into haloes of mass $> M_{\text{min}}$, $F_{\text{coll}}(> M_{\text{min}}, z)$. The shaded region shows the collapsed fraction in CDM, with a range of low-mass cutoffs corresponding to virial temperatures $300 \text{ K} < T_{\text{vir}} < 10^4 \text{ K}$. The other curves correspond to WDM particle masses of $m_x = 3.0, 2.5, 2.0, 1.5,$ and 1.0 keV (top to bottom). This figure is analogous to Fig. 2 in Mesinger et al. (2005), serving to motivate equation (4). The fractions $F_{\text{coll}}(> M_{\text{min}}, z)$, computed according to the full random walk procedure used in Mesinger et al. (2005), are shown as solid black curves, while the approximation of a sharp cutoff at M_{WDM} (eq.4) corresponds to the red dashed curves.

The comoving star formation rate as a function of redshift is assumed to be proportional to dF_{coll}/dt :

$$SFR(z) \propto \frac{dF_{\text{coll}}}{dt}. \quad (10)$$

The proportionality constant is irrelevant for our analysis, as it is subsumed in our normalization procedure below.

³ <http://camb.info/>

⁴ Using $M_{\text{gal}}(z)$ from Sobacchi & Mesinger (2013a) for WDM models is not entirely self-consistent, since UVB feedback effects should be smaller in WDM models. However, this effect is smaller than other astrophysical uncertainties. Most importantly, our conclusions are driven by models which at high redshift have $M_{\text{WDM}} > M_{\text{gal}}$, and are therefore insensitive to the actual choice of M_{gal} (see Fig. 2).

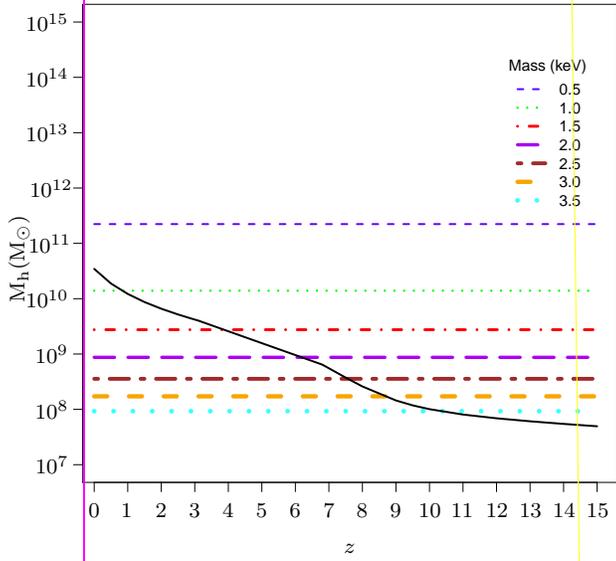
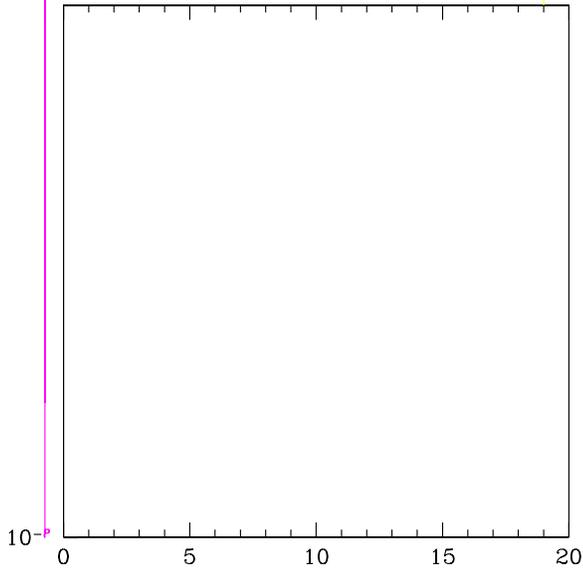


Figure 1. The minimum halo masses capable of hosting star-forming galaxies. The solid black line corresponds to the astrophysical limit, M_{gal} , from Sobacchi & Mesinger (2013a). The horizontal lines correspond to the cosmological cutoffs in WDM models, M_{WDM} ($m_x = 0.5, 1, 1.5, 2, 2.5, 3, 3.5$ keV from top to bottom, respectively).



time interval in the observer rest frame⁷, and $\mathcal{I}(z)$ is the integral over the GRB luminosity function (LF), $p(L)$,

$$\mathcal{I}(z) = \int_{\log L_{\text{lim}}(z)}^{\infty} p(L) d \log L. \quad (14)$$

To remove the dependence on proportionality constants, we construct cumulative distribution functions (CDFs) of GRBs over the redshift range $z_i < z < z_{\text{max}}$:

$$N(< z | z_{\text{max}}) = \frac{N(z_i, z)}{N(z_i, z_{\text{max}})}. \quad (15)$$

The expected number of observed GRBs in a given redshift interval, for each combination $\theta \equiv \{m_x, \alpha\}$, can be written as

$$N(z_1, z_2; \theta) = \zeta_{0;\theta} \int_{z_1}^{z_2} (1+z)^\alpha \text{SFR}(z; m_x) \frac{\Delta t}{1+z} \frac{dV}{dz} \mathcal{I}(z), \quad (16)$$

where $\zeta_{0;\theta}$ is normalized so that each model recovers the observed GRB rate, $N_{\text{obs}}(z_1, z_2)$, at $z \approx 4$,

$$\zeta_{0;\theta} = \frac{N_{\text{obs}}(3, 4)}{\int_3^4 (1+z)^\alpha \text{SFR}(z; m_x) \frac{\Delta t}{1+z} \frac{dV}{dz} \mathcal{I}(z)}. \quad (17)$$

$N_{\text{obs}}(3, 4)$ is equal to 24 (7) for S1 (S2) samples respectively. Note that normalizing at a relatively high- z ($z=3-4$) is indeed a conservative choice. If we had normalized at lower redshifts, say $z \sim 1-2$, the absolute number of GRBs predicted by $m_x = 0.5$ keV and CDM, between $z \sim 3-4$, would already diverge by a factor of ~ 2 for α values in the range $0-2$. In addition to being conservative, normalizing at high-redshifts allows us to have a relatively short lever arm over which our simple scaling, $\text{SFR} \propto dF_{\text{coll}}/dt$, is presumed to be accurate. At lower redshifts, mergers and AGN feedback (missing from our model) are expected to be important in determining the SFR. Hence to normalize our CDFs, we chose the largest redshift at which our sample is reasonably large (see below). It is important to note however that our main results are based on the $z > 4$ CDFs, which unlike the predictions for the absolute numbers of bursts, *do not depend on our choice of normalization*.

4 GRB SAMPLE

Our LGRB data is taken from Robertson & Ellis (2012), corresponding to a compilation from the samples presented in Butler et al. (2007); Perley et al. (2009); Butler et al. (2010); Sakamoto et al. (2011); Greiner et al. (2011), and Krühler et al. (2011). It includes only GRBs before the end of the Second Swift BAT GRB Catalog, and is comprised of 152 long GRBs with redshift measurements. It is important to consider the completeness of the sample. Several efforts have been made to construct a redshift-complete GRB sample (e.g., Greiner et al. 2011; Salvaterra et al. 2012). However, to do so, many GRBs with measured redshifts are excluded. Such requirements are even more severe for high- z bursts, which makes them of little use for our purposes. To explore the dependence of our results on a possible bias in

⁷ Which in our case, corresponds to 5 yrs of observation time by *Swift*. Again, the value is not relevant for our purposes, since it is subsumed in the normalization.

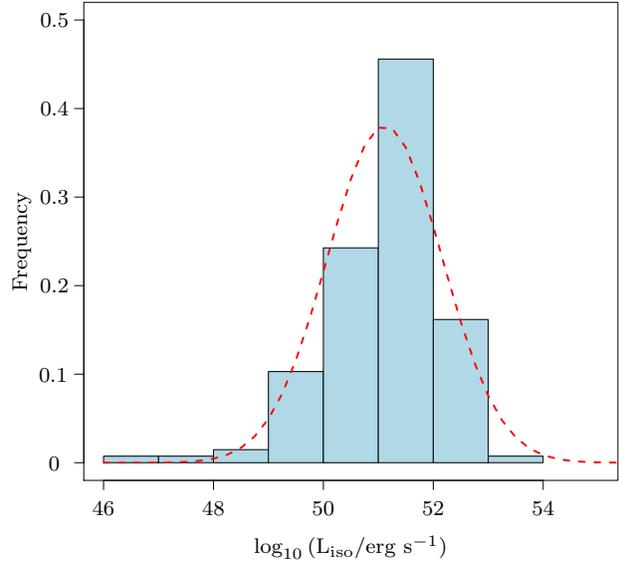


Figure 3. Frequency (i.e. fraction in bin) of GRB luminosities for the $z < 4$ subsample used to construct the LF used for S1 (see text for details). The red dashed line represents the best-fit LF.

the GRB redshift distribution, we construct two samples: (S1) we use a luminosity function based on the $z < 4$ subsample (consisting of 136 GRBs); and (S2) we use a subsample with isotropic-equivalent luminosities bright enough to be observable up to high redshifts (comprised of 38 bursts). The two samples are summarized in table 1. Since there is a degeneracy between a biased SFR-GRB relation and a redshift-dependent LF, we implicitly assume that any unknown bias will be subsumed in the value of the α parameter.

4.1 Luminosity function sample (S1)

The number of GRBs detectable by any given instrument depends on the specific flux sensitivity threshold and the intrinsic isotropic LF of the GRBs. In figure 3, we show the distribution of log-luminosities for $z < 4$ GRBs, which can be well described by a normal distribution

$$p(L) = p_* e^{-(L-L^*)^2} / 2\sigma_L^2. \quad (18)$$

The values $L = \log L_{\text{iso}}/\text{erg s}^{-1}$; $L^* = \log 10^{51.16}$, $\sigma_L = 1.06$ and $p_* = 1.26$ are estimated by maximum likelihood optimization. The luminosity threshold is given by

$$L_{\text{lim}} = 4\pi d_L^2 F_{\text{lim}}, \quad (19)$$

where d_L is the luminosity distance. Consistent with previous works (e.g. Li 2008), we set a bolometric energy flux limit $F_{\text{lim}} = 1.15 \times 10^{-8} \text{erg cm}^{-2} \text{s}^{-1}$ for *Swift* by using the smallest luminosity of the sample. Due to Malmquist bias, our fitted LF is likely biased towards high luminosities. To minimize this problem and increase the sample completeness, we fit our LF using only $z \leq 4$ GRBs, which we show in Fig. 3. We reiterate that the Malmquist bias only serves to make our results even more conservative by predicting a flatter redshift distribution of GRBs.

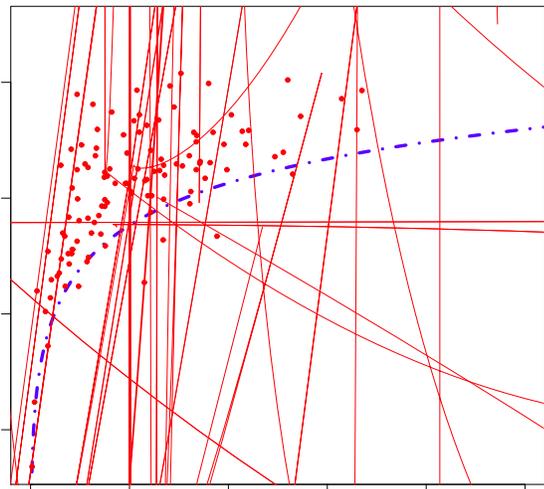


Table 2. Absolute number of GRBs per redshift interval predicted by each model for S1C1 sample.

Model	N(4,6)			N(6,8)			N(8,10)		
	$\alpha = 0$	$\alpha = 1$	$\alpha = 2$	$\alpha = 0$	$\alpha = 1$	$\alpha = 2$	$\alpha = 0$	$\alpha = 1$	$\alpha = 2$
$m_x = 0.5$ keV	3.18	3.94	4.88	0.01	0.02	0.04	1.0×10^{-5}	2.2×10^{-5}	4.6×10^{-5}
$m_x = 1.0$ keV	9.34	11.82	15.00	0.42	0.71	1.21	0.01	0.02	0.04
$m_x = 1.5$ keV	14.84	19.10	24.67	1.51	2.58	4.42	0.08	0.17	0.36
$m_x = 2.0$ keV	14.31	18.36	23.65	2.31	4.01	6.96	0.20	0.43	0.93
$m_x = 2.5$ keV	14.44	18.54	23.90	2.09	3.65	6.38	0.39	0.83	1.80
$m_x = 3.0$ keV	14.51	18.62	24.01	2.04	3.54	6.16	0.46	1.00	2.20
$m_x = 3.5$ keV	14.56	18.69	24.09	2.04	3.55	6.18	0.44	0.97	2.13
CDM	14.56	18.69	24.10	2.04	3.55	6.18	0.45	1.00	2.19
<i>Swift</i>	11	11	11	3	3	3	2	2	2

Table 3. Absolute number of GRBs per redshift bin predicted by each model for S2C1 sample.

Model	N(4,6)			N(6,8)			N(8,10)		
	$\alpha = 0$	$\alpha = 1$	$\alpha = 2$	$\alpha = 0$	$\alpha = 1$	$\alpha = 2$	$\alpha = 0$	$\alpha = 1$	$\alpha = 2$
$m_x = 0.5$ keV	1.33	1.65	2.05	0.01	0.02	0.03	1.5×10^{-5}	3.1×10^{-5}	6.6×10^{-5}
$m_x = 1.0$ keV	4.10	5.23	6.70	0.34	0.59	1.00	0.01	0.03	0.06
$m_x = 1.5$ keV	6.75	8.78	11.46	1.27	2.19	3.76	0.12	0.26	0.55
$m_x = 2.0$ keV	6.47	8.40	10.94	2.02	3.52	6.14	0.31	0.67	1.44
$m_x = 2.5$ keV	6.54	8.49	11.07	1.86	3.27	5.77	0.60	1.30	2.81
$m_x = 3.0$ keV	6.57	8.53	11.12	1.79	3.13	5.49	0.74	1.63	3.58
$m_x = 3.5$ keV	6.59	8.56	11.16	1.80	3.14	5.51	0.72	1.58	3.48
CDM	6.60	8.56	11.16	1.80	3.15	5.51	0.74	1.63	3.58
<i>Swift</i>	6	6	6	3	3	3	2	2	2

5.3 Constraints from the redshift distribution of $z > 4$ bursts

To quantify how consistent are these CDFs with the observed distribution from *Swift*, we make use of two statistics: (i) the one-sample Kolmogorov-Smirnov (K-S) test; and (ii) a maximum likelihood estimation (MLE). Both tests are described in detail in appendix A.

The K-S test provides a simple estimate of the probability the observed distribution was drawn from the underlying theoretical one. We compute this probability, for fixed α first, for our models S1C0, S1C1, S2C0 and S2C1. Consistent with the more qualitative analysis from the previous section, models with $m_x \lesssim 1.0$ keV are ruled out at 90% CL assuming $-1 \leq \alpha \leq 1$. For $\alpha = 0$ (S1C0 and S2C0), the limits are even more restrictive and models with $m_x \lesssim 1.5$ keV are ruled out at 90% CL for both samples.

So far, we have analyzed each model individually in order to quantify a lower limit on m_x , given a single value of α . Using a χ^2 MLE (see appendix A) allows us to compute posterior probabilities given conservative priors on α . Thus we are able to construct confidence limits in the two-dimensional, (m_x, α) parameter space. The results for cases S1C2, S2C2 are shown in figure 7 at 68%, 95%, 99% CL. Both samples show the same qualitative trends, with the data preferring higher values of m_x and CDM. Marginalizing the likelihood over $-3 \leq \alpha \leq 12$, with a flat prior, shows that models with $m_x \leq 1.6 - 1.8$ keV are ruled out at 95% CL for S1C2 and S2C2 respectively.

6 FUTURE CONSTRAINTS

In the previous section, we have quantified the constraints on WDM particle masses using current *Swift* GRB obser-

vations. We obtain constraints of $m_x \gtrsim 1.6 - 1.8$ keV. We now ask how much could these constraints could improve with a larger GRB sample, available from future missions? As a reference, we use the Sino-French space-based multi-band astronomical variable objects monitor (SVOM)⁹ mission. The SVOM has been designed to optimize the synergy between space and ground instruments. It is forecast to observe $\sim 70 - 90$ GRBs yr^{-1} and $\sim 2 - 6$ GRB yr^{-1} at $z \geq 6$ (see e.g., Salvaterra et al. 2008).

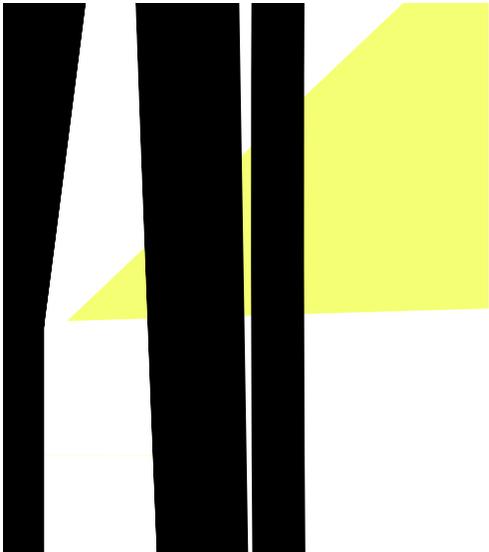
We first construct a mock GRB dataset of 450 bursts with redshifts obtained by sampling the CDM, $\alpha = 0$ PDF given by equation (13). This sample size represents an optimistic prediction for 5 yrs of SVOM observations¹⁰ (see e.g., Salvaterra et al. 2008). We then perform the MLE analysis detailed above on this mock dataset at $z > 4$. The resulting confidence limits are presented in Fig. 8.

This figure shows that ~ 5 yrs of SVOM observations would be sufficient to rule out $m_x \leq 2.3$ keV models (from our fiducial CDM, $\alpha = 0$ model) at 95% CL, when marginalized over α . This is a modest improvement over our current constraints using *Swift* observations. As already foreshadowed by figures 1 and 2, as well as the associated discussion, it is increasingly difficult to push constraints beyond $m_x > 2$ keV. On the other hand, the α constraint improves dramatically due to having enough high- z bursts to beat the Poisson errors. We caution that the relative narrowness around $\alpha = 0$ of the contours in Fig. 8 is also partially due to our choice of (CDM, $\alpha = 0$) as the template for the mock observation.

⁹ <http://www.svom.fr>

¹⁰ The highest redshift in our mock sample is $z_{\text{max}} = 11.6$, being the only event at $z > 10$.

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APPENDIX A: STATISTICS

Kolmogorov-Smirnov one-sample test

A straightforward way to compare the data and WDM models is to perform a one-sample Kolmogorov-Smirnov (K-S) test. The null hypothesis that the observed GRB redshifts are consistent with a model distribution can be evaluated by estimating a *p-value*, which corresponds to one minus the probability that the null hypothesis can be rejected. The K-S test consists in comparing the statistical parameter

$$D = \sup|F(z) - G(z)|, \quad (\text{A1})$$

where $F(z)$ and $G(z)$ are the CDF for the theoretical and observed sample and \sup is the supremum of a totally or partially ordered set. We estimate the *p-value* via nonparametric bootstrap, which consists of running Monte Carlo realizations of the observed CDF using a random-selection-with-replacement procedure estimated from the data. This provides a histogram of the statistic D , from which a valid goodness-of-fit probability can be evaluated. The probability distribution function for each model is determined by equation (16).

Maximum-likelihood method

More formally, we can estimate the probability of parameters $\{\alpha, m_x\}$ given the observed data using a Bayesian technique. Assuming that our data is described by the probability density function $f(x; \boldsymbol{\theta})$, where x is a variable and $\boldsymbol{\theta} \equiv \{\alpha, m_x\}$. We want to estimate $\boldsymbol{\theta}$, assuming the data are independent. So the likelihood will be given by

$$\mathcal{L}(\boldsymbol{\theta}|x_i) \propto \prod_{i=1}^N f(x_i|\boldsymbol{\theta}). \quad (\text{A2})$$

Given the small number of observed bursts per redshift bin, $\Delta z = 1.5$, we use a Poisson error statistics¹¹ (see e.g., Campanelli et al. 2012 for a similar procedure applied to galaxy cluster number count). Therefore, the likelihood function can be computed as

$$\mathcal{L}(\boldsymbol{\theta}|\kappa_i) \propto \prod_{i=1}^5 \frac{\Upsilon_i^{\kappa_i} e^{-\Upsilon_i}}{\kappa_i!}, \quad (\text{A3})$$

where $\Upsilon_i \equiv N(z_i, z_{i+1}; \boldsymbol{\theta})$ and $\kappa_i \equiv N_{\text{obs}}(z_i, z_{i+1})$. Thus, the χ^2 statistics can be written as

$$\begin{aligned} \chi^2(\boldsymbol{\theta}) &= -2 \ln \mathcal{L}, \\ &= 2 \sum_{i=1}^5 \Upsilon_i - \kappa_i (1 + \ln \Upsilon_i - \ln \kappa_i). \end{aligned} \quad (\text{A4})$$

¹¹ The value $\Delta z = 1.5$ is chosen to ensure at least 1 burst per bin.